

# Eigenmodes of Susceptibility Matrix (帯磁率行列の固有モード解析)

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**In collaboration with Prof. Iba (ISM)**

# Outline

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- Correlation Matrix Spectral Analysis
- New techniques for numerical diagonalization of correlation matrix
  - Dual Trick
  - (On-line Method)
- Applications to spin models
  - three-dimensional fully frustrated Ising model
  - short-range Ising spin glass model

# Eigenmode Analysis of Susceptibility Matrix

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Susceptibility Matrix ( $N \times N$  dense matrix)

$$\boxed{C_{ij} = \langle S_i S_j \rangle} = \begin{cases} \delta_{ij} & (T \rightarrow \infty) \\ +1 & (T \rightarrow 0 \text{ in non-frustrated ferromagnet}) \end{cases}$$

Eigenmode associated with  $O(N)$  eigen value : Long range order

$$\chi_{SG}^J \equiv \frac{1}{N} \sum_{ij} (\langle S_i S_j \rangle - \langle S_i \rangle \langle S_j \rangle)^2 = \frac{1}{N} \sum_{ij} C_{ij}^2 = \text{Tr } C^2 / N \quad \text{for } T > T_c$$

- Analytical argument of Griffiths phase in random spin systems :  
A.J.Bray and M.A.Moore J. Phys. **C15** (1982) L765.  
A.J.Bray, Phys. Rev. Lett., **59** (1987) 586
- Numerical studies of the Griffiths phase: Nemoto-Yamada, unpublished
- Dynamical correlation matrix:  
H. Takano and S. Miyashita, J. Phys. Soc. Jpn., **64** (1995) 3688.

# Correlation-Matrix Analysis in the Spin Glass Phase (1)

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- J. Sinova, G. Canright and A. H. MacDonald, Phys. Rev. Lett. **85** (2000) 2609.
- J. Sinova, G. Canright, H. E. Castillo and A. H. MacDonald, Phys. Rev. B **85**, (2001) 104427. (J. Sinova and G. Canright, cond-mat/0103071.)  
Random Field Ising Model

Ordering appears in the spectrum of the correlation matrix  
**AS AN EXTENSIVE EIGENVALUE.**

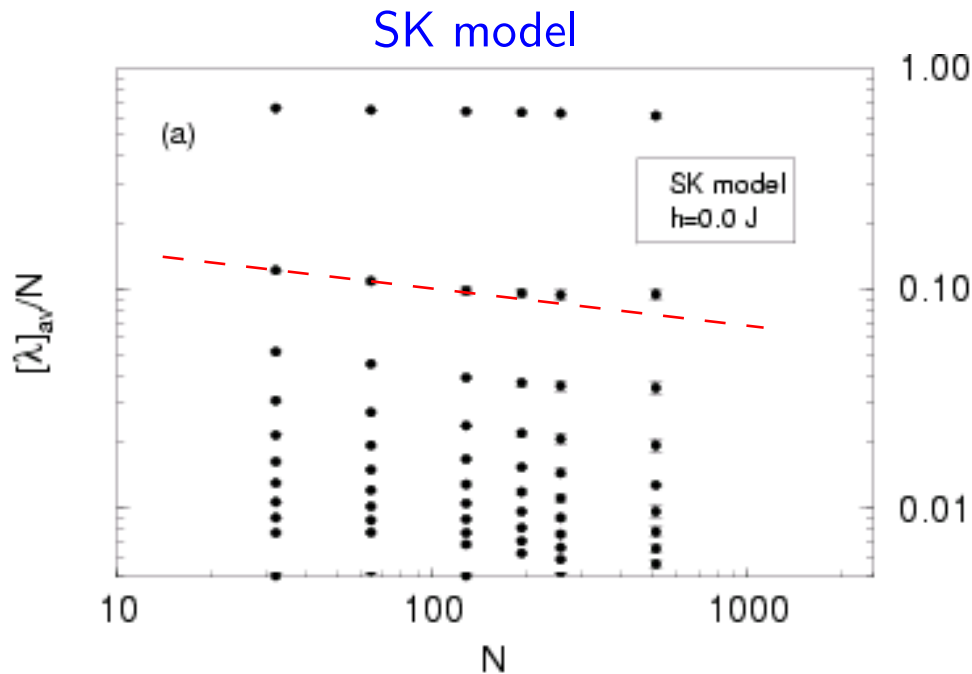
There is **RSB**  $\longrightarrow$  There must be **MORE THAN ONE EXTENSIVE EIGENVALUE** in the spectrum of  $C_{ij}$ .

Absence of RSB  $\longrightarrow$  There can only be one such extensive eigenvalue.

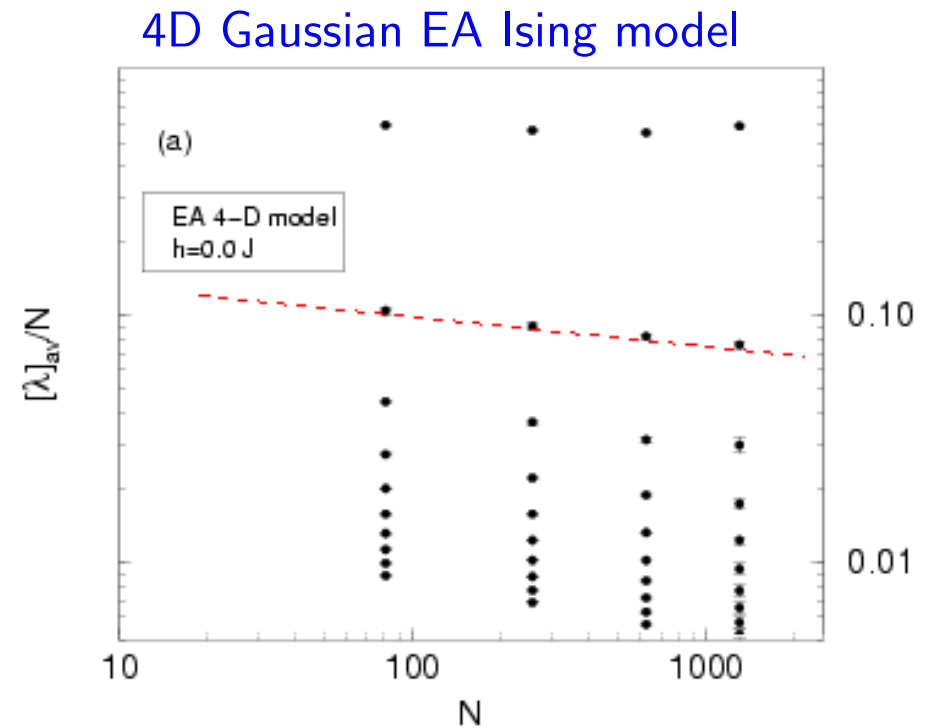
**# of Extensive Eigenvalues  $\rightarrow$  # of pure states.**

## Correlation-Matrix Analysis in the Spin Glass Phase (2)

Finite size scaling of the disorder average of the ten largest eigenvalues of  $C_{ij}$



**Multiple extensive eigenvalues** in the RSB phase, as expected.



**A single extensive eigenvalue**, consistent with the droplet theory.

# Conventional Method for the Eigenmode Analysis

STEP 1: Evaluate susceptibility-matrix using Monte Carlo simulations.

$$C_{ij} = \langle (S_i - \bar{S}_i)(S_j - \bar{S}_j) \rangle = \frac{1}{M} \sum_{\mu}^M (S_i^{\mu} - \bar{S}_i)(S_j^{\mu} - \bar{S}_j) \quad (1)$$

where  $S_i = \frac{1}{M} \sum_{\mu} S_i^{\mu}$  and  $M$  is the total number of MC steps.

STEP 2: Diagonalize the matrix  $C_{ij}$

$$\sum_j^N C_{ij} e_j = \lambda e_i. \quad (2)$$

Eigenmode Analysis in Physics

Principal Component Analysis (PCA) in  
Statistics

- $O(N^2)$  memory to storage the matrix element of  $C_{ij}$  with  $N \times N$ .
- $O(N^3)$  time complexity to diagonalize  $C_{ij}$  by Householder method.

$$\implies \begin{cases} \text{Nemoto et al} & N \leq 28^2 = 784 \\ \text{Sinova et al} & N \leq 6^4 = 1296 \rightarrow 13\text{MB} \end{cases}$$

# New Methods for Eigenmode Analysis

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Don't keep the matrix elements of  $C_{ij}$ .

## Proposed Methods

### 1. **Dual Trick**

- Memory :  $O(N)$
- Time:  $O(N)$
- Drawback: diagonalization of  $M \times M$  matrix

### 2. **One-line Trick**

- Memory:  $O(N \times K)$  ,  $K$  is a number of our desired eigenmodes.
- Time:  $O(N \times K^2)$
- Drawback: Need iteration for many times.

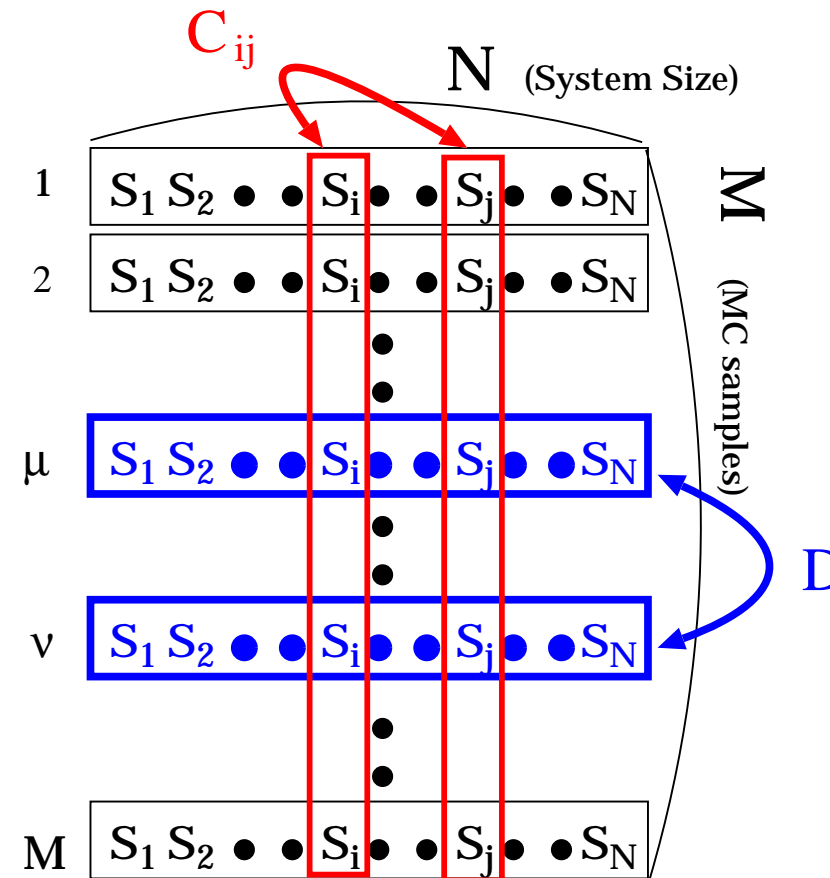
# Dual Trick Method

Step 1: Evaluate **Hamming-distance Matrix**  $D^{\mu\nu}$  between data obtained by MC simulation

$$D^{\mu\nu} = \frac{1}{N} \sum_i^N (S_i^\mu - \bar{S}_i)(S_i^\nu - \bar{S}_i).$$

Step 2: Diagonalize the Matrix  $D^{\nu\mu}$  (**Dual Problem**)

$$\sum_\nu^M D^{\mu\nu} \hat{e}_\nu = \lambda \hat{e}_\mu.$$





## Dual to Original Problem

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- Dual relation

$$\text{primary} \quad \sum_j^N C_{ij} e_j = \lambda e_i. \quad (N \times N)$$

$$\text{dual} \quad \sum_\nu^M D^{\mu\nu} \hat{e}_\nu = \lambda \hat{e}_\mu. \quad (M \times M)$$

$$\text{DATA} \quad X_i^\mu = (S_i^\mu - \langle S_i \rangle) \quad (M \times N)$$

$$e_i = \sum_{\mu=1}^M X_i^\mu \hat{e}^\mu$$

Eigenvector with the eigenvalue  $\lambda/M$   
of the correlation matrix  $C_{ij}$

- Paradox: Missing rank ?

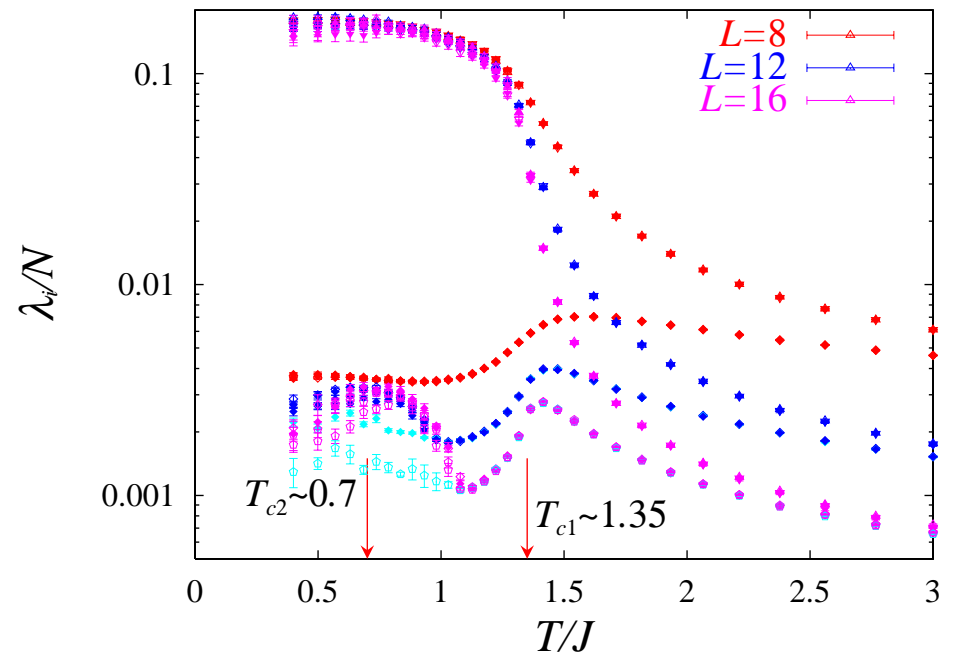
$$[\text{rank}] = \min(N, M)$$

# Application to 3d Fully Frustrated Ising Model

## 3d Fully Frustrated Ising Model

using **On-line trick method**

- Multiple pure states
  - Phase Transition 1:  
 $T_{c1}/J \sim 1.35$   
16-fold states
  - Phase Transition 2:  
 $T_{c2}/J \sim 0.7$   
24-fold states



Temperature dependence of the eigenvalues.

# Application to Ising Spin Glass Model

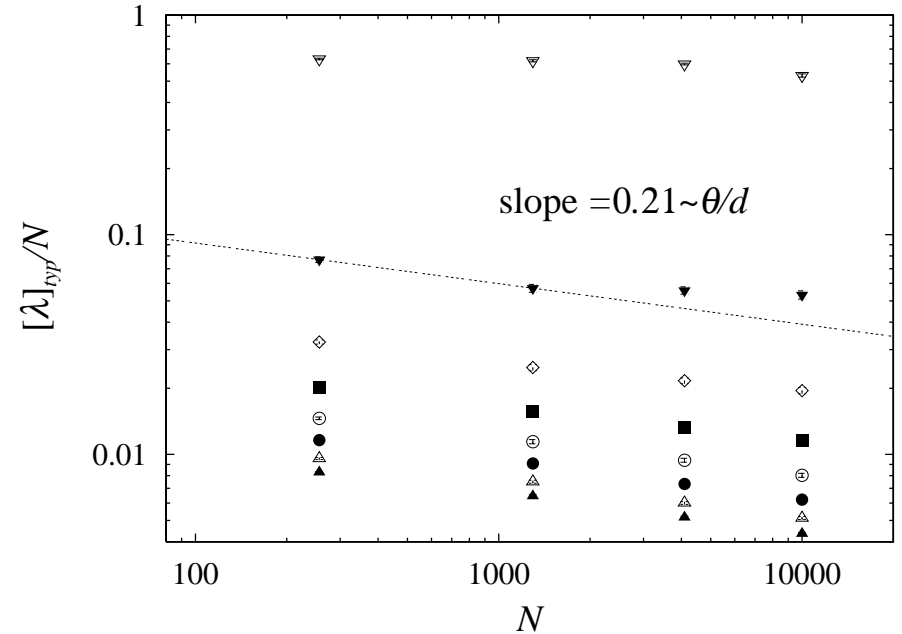
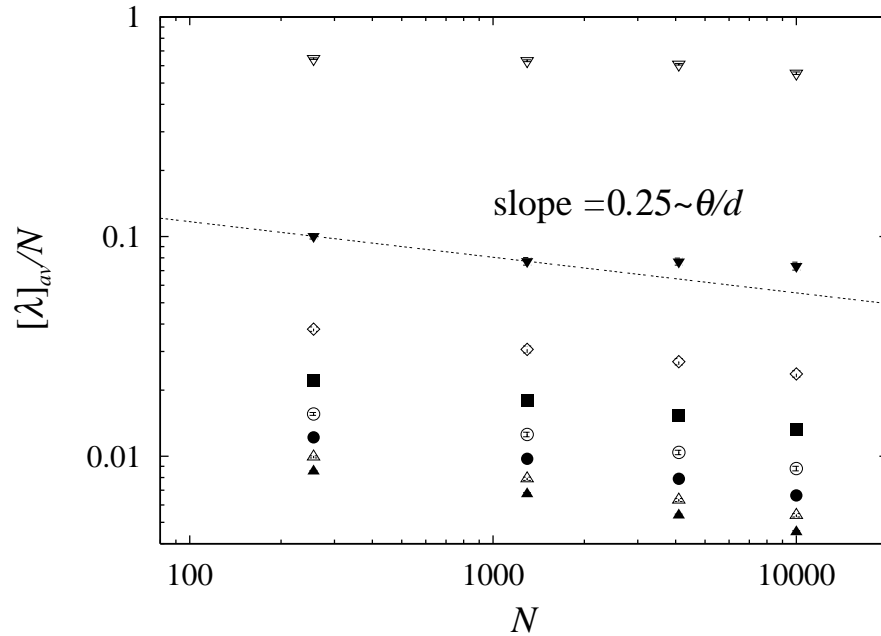
$4d \pm J$  Ising EA Model using the dual trick method

System Size:  $N \leq 10^4 = 10,000$  (800MB in conventional method)

Temperature:  $T/J = 1.0 \sim 0.5T_c$

Average

Typical ( $[\lambda]_{\text{type}} \equiv \exp[\ln \lambda_i]_{\text{av}}$ )



The second eigenvalue is found to be **extensive** for  $N > 6^4$ .

$\Rightarrow$  Multiple Pure States

# Summary

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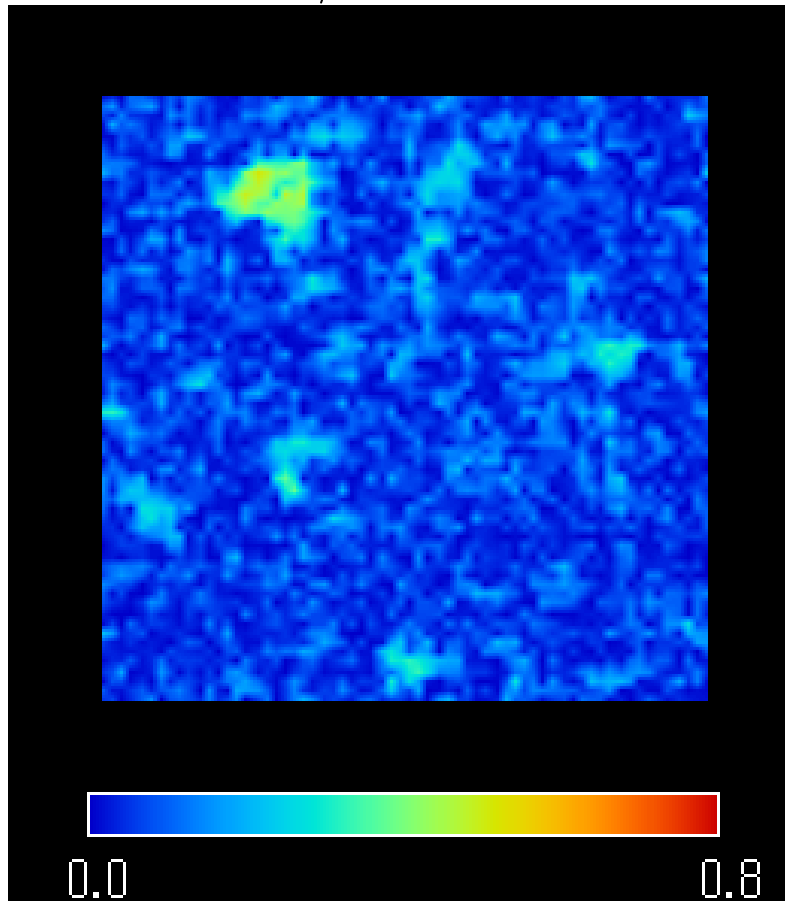
- New technique for Eigenmode Analysis to correlation-matrix
  - Dual Trick  
from Hamming-distance matrix to spin-correlation matrix
  - (On-line Trick)
- Application to  $4d \pm J$  Ising EA model

The second eigenvalue is also extensive at low temperatures.  
 $\implies$  Multiple pure states
- Future problems
  - Eigenvector  $\longrightarrow$  Ordering pattern, Site magnetization
    - \* Direct observation of Low-energy excitations like droplet

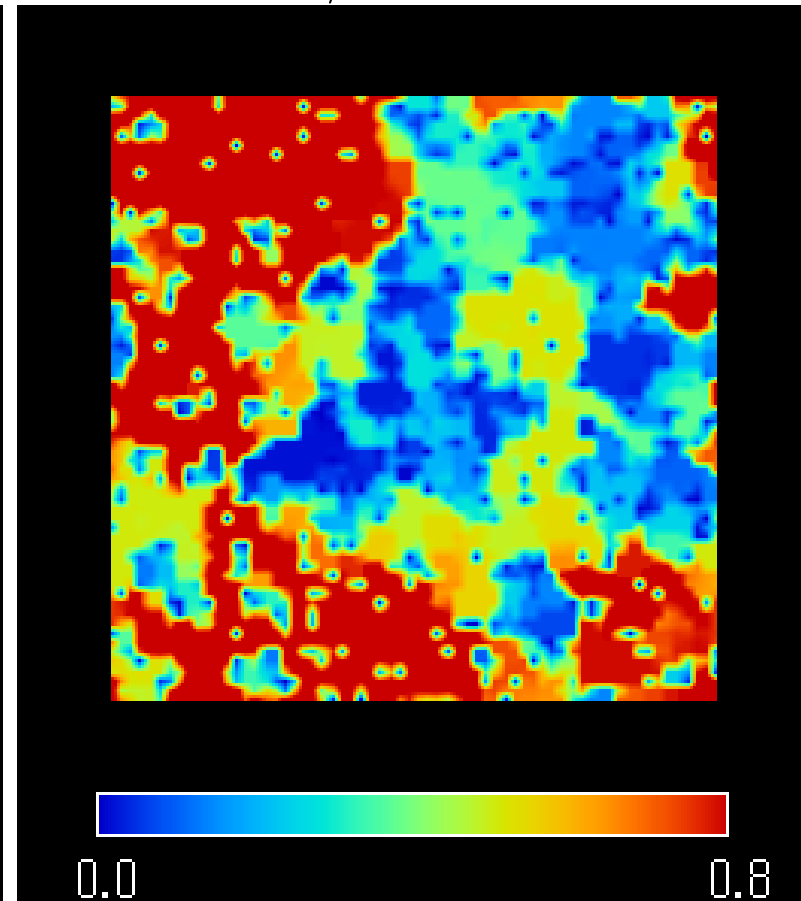
# Application to $2d$ Ising Spin Glass Model

The first eigenvector of a sample with  $N = 64^2$

$T/J = 1.8$

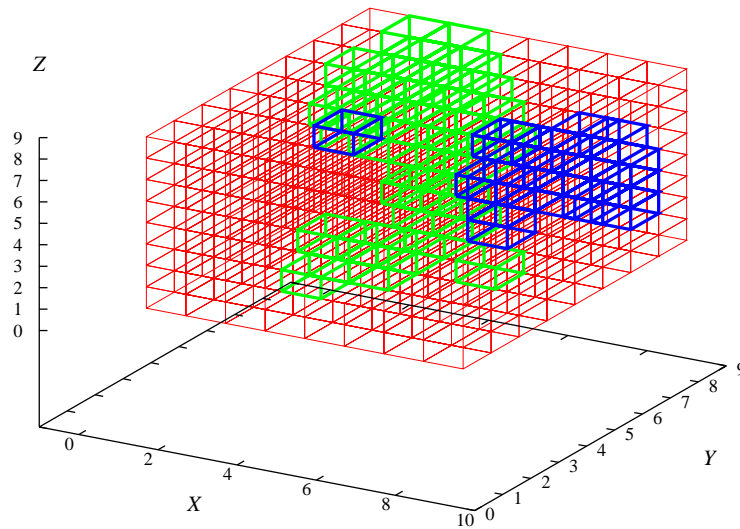


$T/J = 0.4$



# Application to $3d \pm$ Ising Spin Glass Model

Three eigenvectors of a sample at  $T/J = 0.4$



Real space plot of spin sites with large amplitude of the **first**(red), **second**(blue) and **third**(green) eigenvectors.

Overlap between data and eigenvector.

